

Numerical stochastic perturbation theory for twisted reduced principal chiral model

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Motivation & Introduction

The principal chiral model for $SU(N)$

$$Z = \int \prod_{\mu, n} dU_{\mu}(n) \exp \left\{ -bN_c \text{Tr} \left(\delta_{\mu} U(n) \delta_{\mu} U^{\dagger}(n) \right) \right\}$$

Rank of matrix

with

$$\delta U(n) = U(n + \mu) - U(n) \quad \beta = bN_c$$

- Two dimensional $SU(N_c) \times SU(N_c)$ principal chiral model - asymptotically free
- Non-perturbative studies on the lattice

- Finite volume correction

$$\langle \mathcal{O}(L) \rangle = \langle \mathcal{O}^{\infty} \rangle - \frac{F(L)}{L^2} + \mathcal{O}(L^4)$$

- Perturbative calculation on the lattice when $N \rightarrow \infty$ limit hold
- QCD like theory
- Too expensive L^2 $SU(N)$ matrix for one configuration

Motivation & Introduction

Twisted Eguchi-Kawai reduction : from lattice to single site

$$S = -bN_c \sum_{\mu} \text{Tr} \left(V_{\mu} V_{\mu}^{\dagger} \right)$$

with

$$V_{\mu} \equiv \Gamma_{\mu} U \Gamma_{\mu}^{\dagger} - U$$

$$\Gamma_1 \Gamma_2 = \exp \left\{ \frac{2\pi i K}{N_c} \right\} \Gamma_2 \Gamma_1 \quad \Gamma_{\mu} \in SU(N_c)$$

- Volume reduction was held in the large-N limit
- Tuning the twist parameter K
 - K, N co-prime
 - $\frac{K}{N} > \Lambda_{cut}$
- Finite-N corrections \leftrightarrow finite-size effects ($L = N_c$)

Motivation & Introduction

- A method that may be able to deal with the perturbation calculation on the lattice up to high-order terms is Numerical Stochastic Perturbation Theory (NSPT)
 - Large N
 - High-order coefficient
- In QFT the asymptotic nature of an expansion in powers of the coupling constant α suggests a factorial divergence occurs when we include high-order coefficients.

Numerical stochastic perturbation theory

- Monte Carlo simulation: Evaluation of Euclidean path integrals for field theories.

- Partition function $Z = \int D\phi \exp[-S[\phi]]$ $\langle O \rangle = \frac{1}{Z} \int D\phi O[\phi] \exp[-S[\phi]]$

Expectation value of observable = Evaluate multi-dimensional integral

- Probabilistic density function $P[\phi] = \frac{1}{Z} \exp[-S[\phi]]$ $\langle O \rangle = \int D\phi O[\phi] P[\phi]$

- Numerical application of stochastic quantization (G.Parisi,Y-S.Wu,1981)

π : Generated from Gaussian noise

- Langevin equation => molecular dynamics based

- Adding an extra DoF $\phi(x) \rightarrow \phi(x, t)$

$$H = \frac{\pi^2}{2} + S[\phi] \quad P'[\pi, \phi] = \frac{1}{Z'} \exp \left[-\frac{\pi^2}{2} - S[\phi] \right]$$

$$\dot{\pi} = -\frac{\partial S}{\partial \phi} = F, \dot{\phi} = \pi \quad \langle O \rangle = \int D\pi D\phi O[\phi] P'[\phi]$$

- Infinite stochastic time limit

$$\left\langle O \left[\phi_\eta(x_1; t) \dots \phi_\eta(x_n; t) \right] \right\rangle_\eta \xrightarrow{t \rightarrow \infty} \left\langle O \left[\phi(x_1) \dots \phi(x_n) \right] \right\rangle$$

Numerical stochastic perturbation theory

- Evaluate the perturbation expansion of physical quantities

$$S[\phi] = S_0[\phi] + S_{Int}[\phi, g]$$

$$\phi = \sum_{k=0}^{N_{trun}} g^k \phi^{(k)}, \pi = \sum_{k=0}^{N_{trun}} g^k \pi^{(k)}$$

- Hierarchical system of partial differential equations

$$\dot{\pi}^{(k)} = -\frac{\partial S^{(k)}}{\partial \phi} [\phi^{(0)}, \phi^{(1)}, \dots, \phi^{(k-1)}], \quad \dot{\phi}^{(k)} = \pi^{(k)} \quad k = 0, 1, 2, \dots$$

$$\langle 0 \rangle = \sum_{k=0} g^k \langle O^{(k)} \rangle = \lim_{T \rightarrow \infty} \frac{1}{T} \sum_{k=0} g^k O^{(k)} [\phi^{(0)}(t), \phi^{(1)}(t), \dots, \phi^{(k)}(t)]$$

NSPT for matrix model

- Convolution operator \circledast

$$A = \sum_{k=0} g^k A^{(k)} = A^{(0)} + gA^{(1)} + g^2A^{(2)} + \dots \quad B = \sum_{k=0} g^k B^{(k)} = B^{(0)} + gB^{(1)} + g^2B^{(2)} + \dots$$

$$AB = C = \sum_{k=0}^{N_{trun}} g^k C^{(k)} = C^{(0)} + gC^{(1)} + g^2C^{(2)} + \dots \quad C^{(l)} = (A \circledast B)^{(l)} = \sum_{k=0}^l A^{(l-k)} B^{(k)}$$

- Expansion for HMD equation(matrix model)

$$\dot{U}^{(i)} = i(\pi \circledast U)^{(i)}$$

$$\dot{\pi}^{(i)} = F^{(i)}$$

- Expansion for $F^{(i)}$

$$F^{(k)} = iN_c \left(V^{(k)} - \frac{1}{N_c} \text{Tr}[V^{(k)}] \right)$$

$$V^{(k)} = S^{(k)} - S^{(k)\dagger}$$

$$S^{(k)} = 4U^{(k)} + X^{(k)} + (U \circledast X)^{(k)} \quad k = 1 \dots N_{truncated}$$

$$X^{(k)} = \sum_{\mu=1}^2 \left[\Gamma_{\mu}^{\dagger} U^{(k)\dagger} \Gamma_{\mu} + \Gamma_{\mu} U^{(k)\dagger} \Gamma_{\mu}^{\dagger} \right] \quad k = 0 \dots N_{truncated}$$

NSPT for matrix model

- Workflow
 - Test NSPT for few N_c and a proper K
 - Calculate internal energy up to $\mathcal{O}(g^8)$
 - Perform $1/N_c^2$ fitting

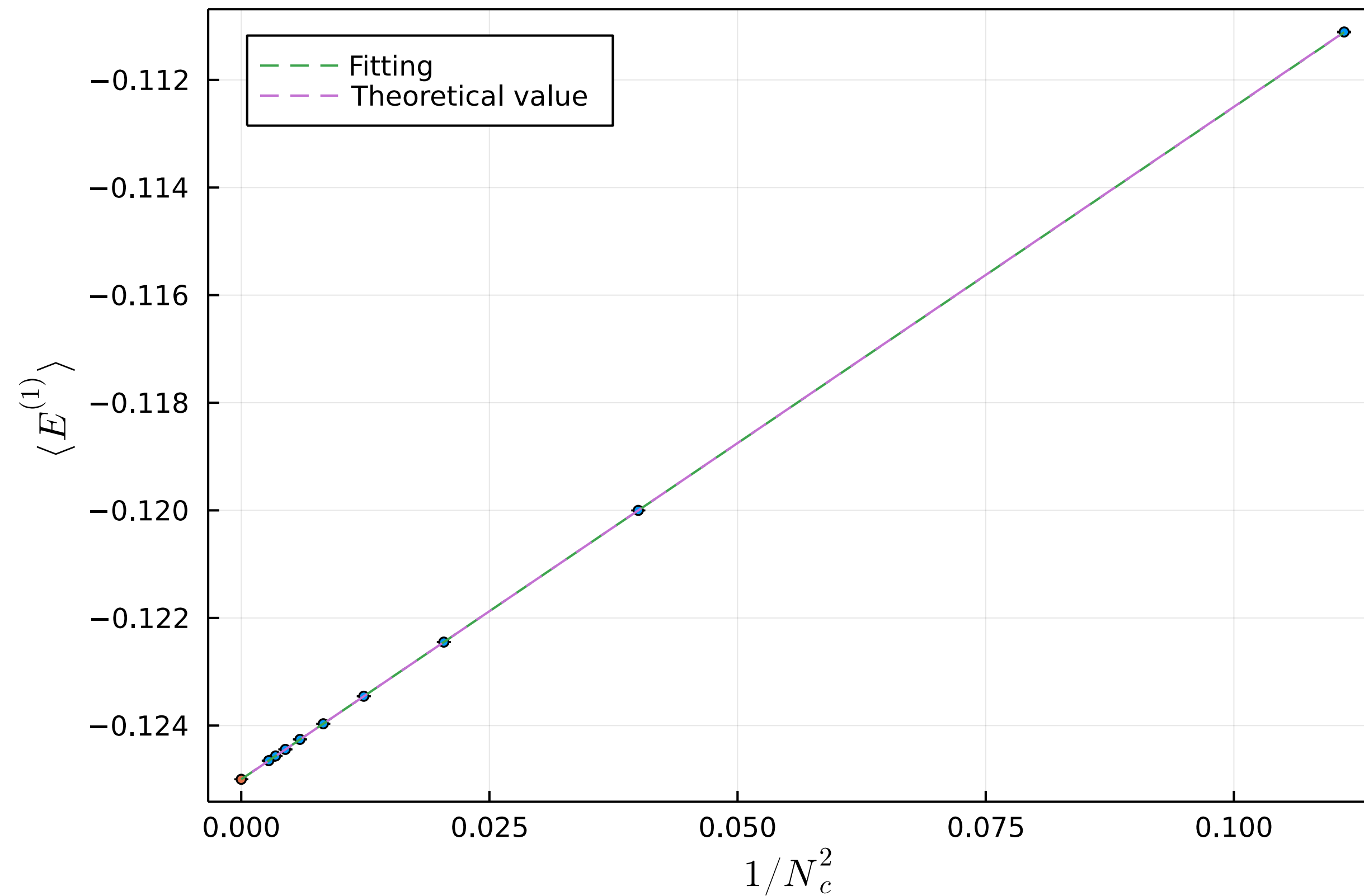
$$E = \frac{1}{2N_c} \sum_{\mu=1,2} \text{Re Tr} [U \Gamma_u U^\dagger \Gamma_u^\dagger]$$

- Few points that need attention
 - The problem with ergodicity in lead order
 - Randomization the trajectory length
 - Partial momentum refreshing scheme $\pi_\mu = c_1 \pi_\mu + \mu + \sqrt{1 - c_1^2} \eta_\mu$
 - Computationally intensive part : convolution operation
 - Naive algorithm $\propto \mathcal{O}(N_{trun}^2)$
 - FFT based convolution $\propto \mathcal{O}(N_{trun} \log N_{trun})$
 - Violation of special unitarity
 - Check the norm of targeted matrix
 - Implement re-unitalization scheme

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Numerical results

1 loop $\mathcal{O}(b)$ case (no volume dependency)

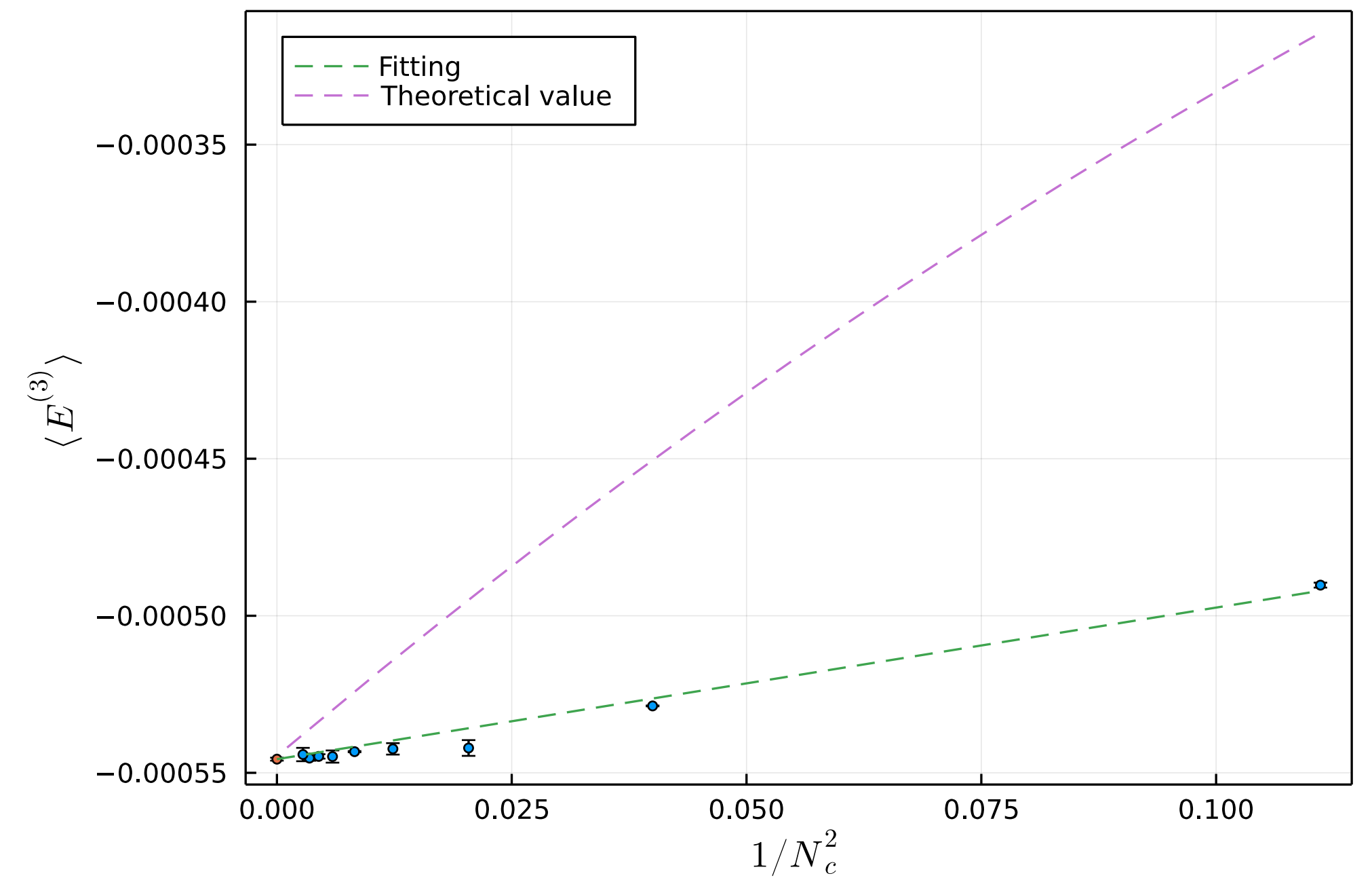
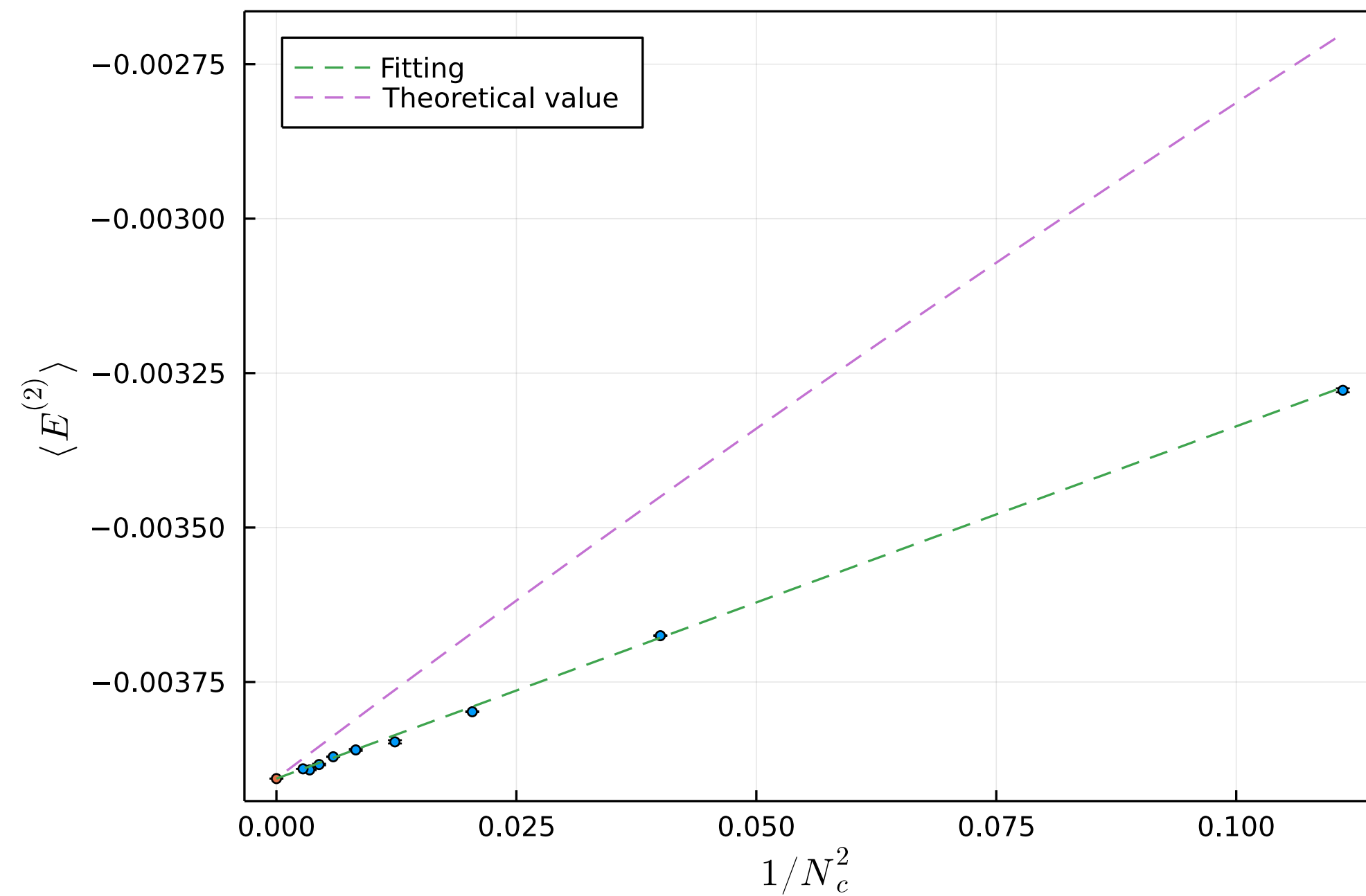


	Theoretical value	NSPT
$E^{(1)}$	$-1/8$	$-0.124999(1)$

Numerical results

2 and 3 $\mathcal{O}(b^{2,3})$ loop case (volume dependency)

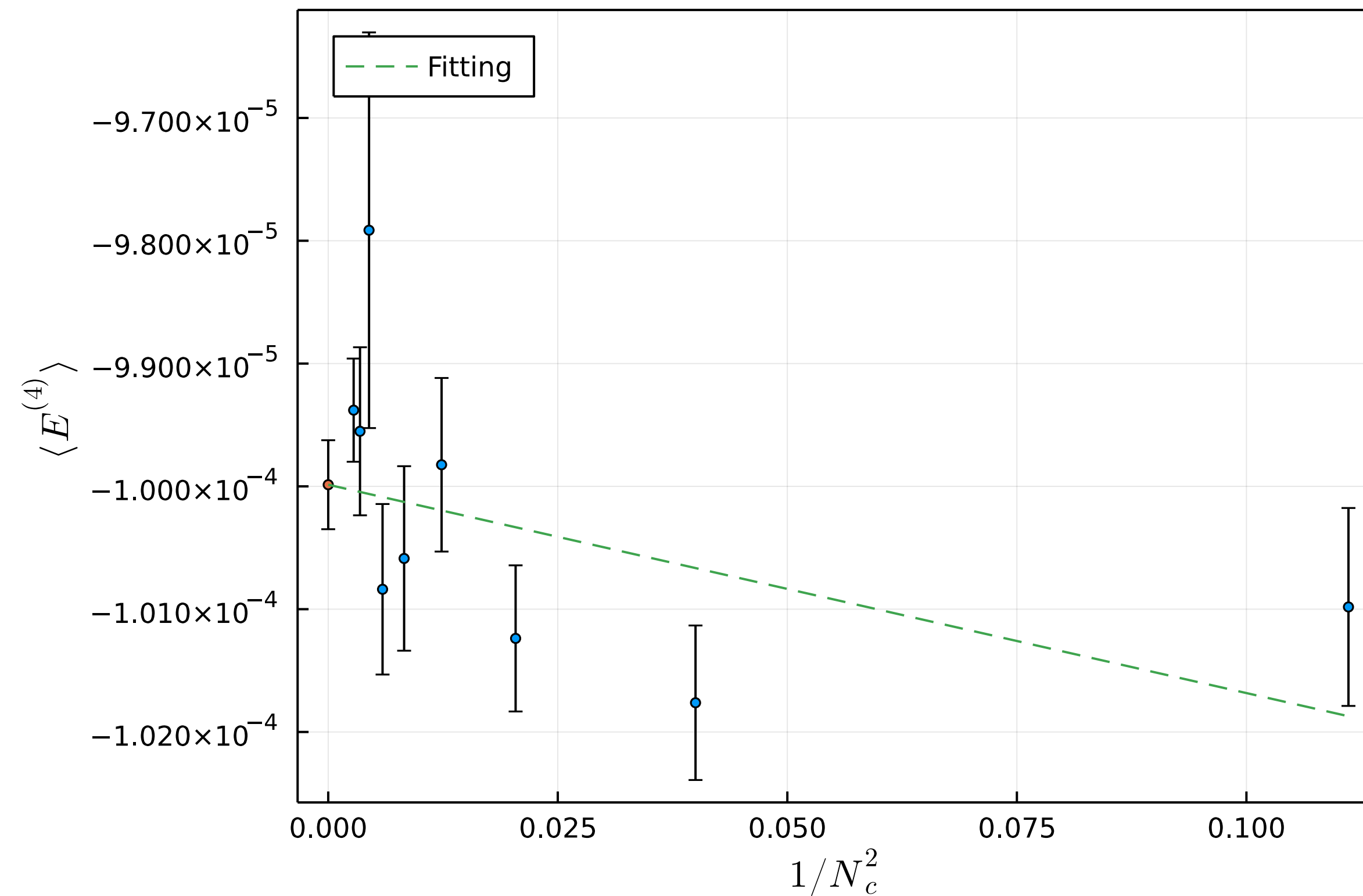
	Theoretical value	NSPT
$E^{(2)}$	-0.00390625	-0.00390637(23)
$E^{(3)}$	-0.000544	-0.000545(4)



Numerical results

4 loop $\mathcal{O}(b^4)$ result

- Our result for four loop case is $E^{(4)} = -0.00009998(36)$

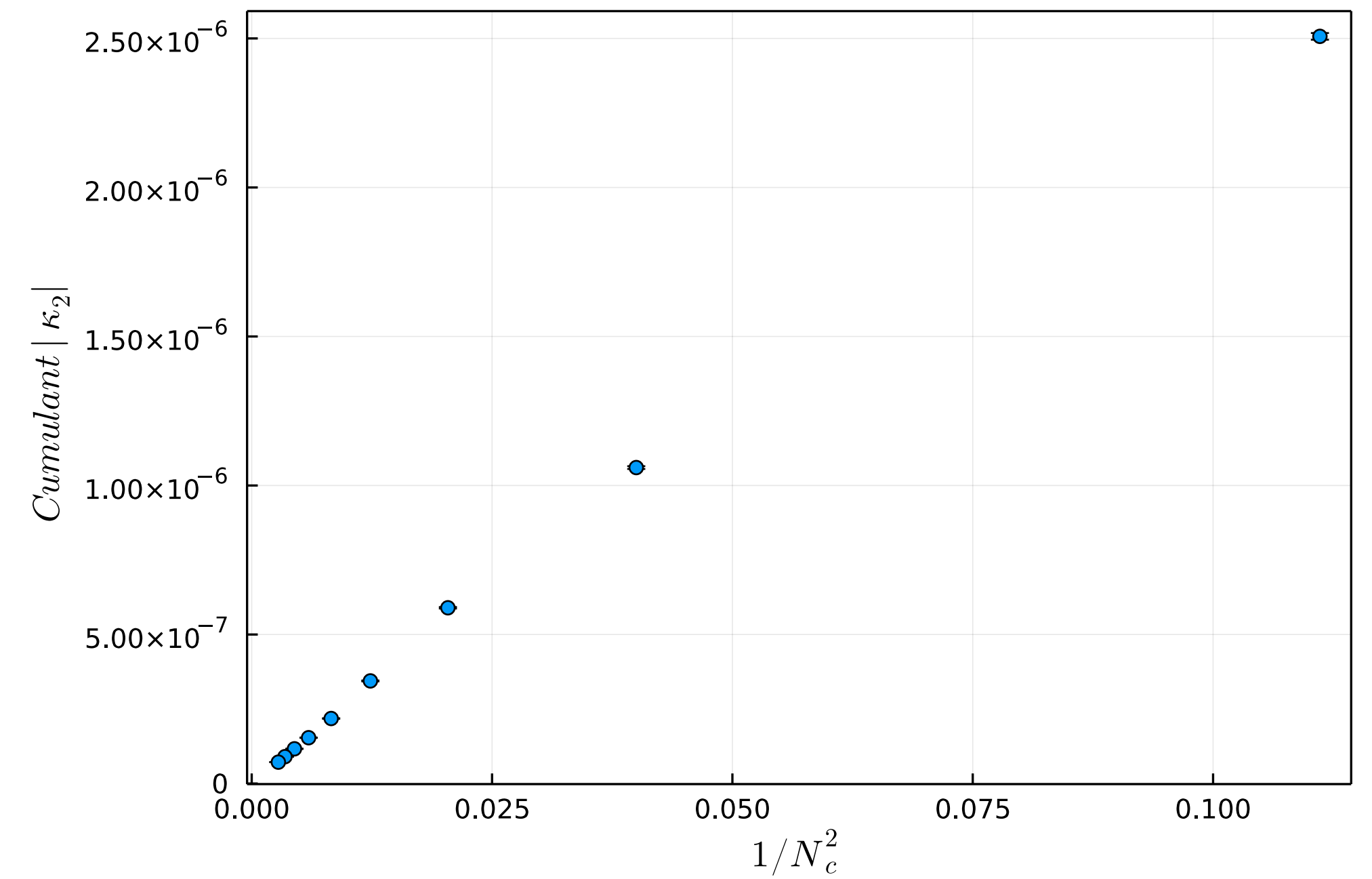
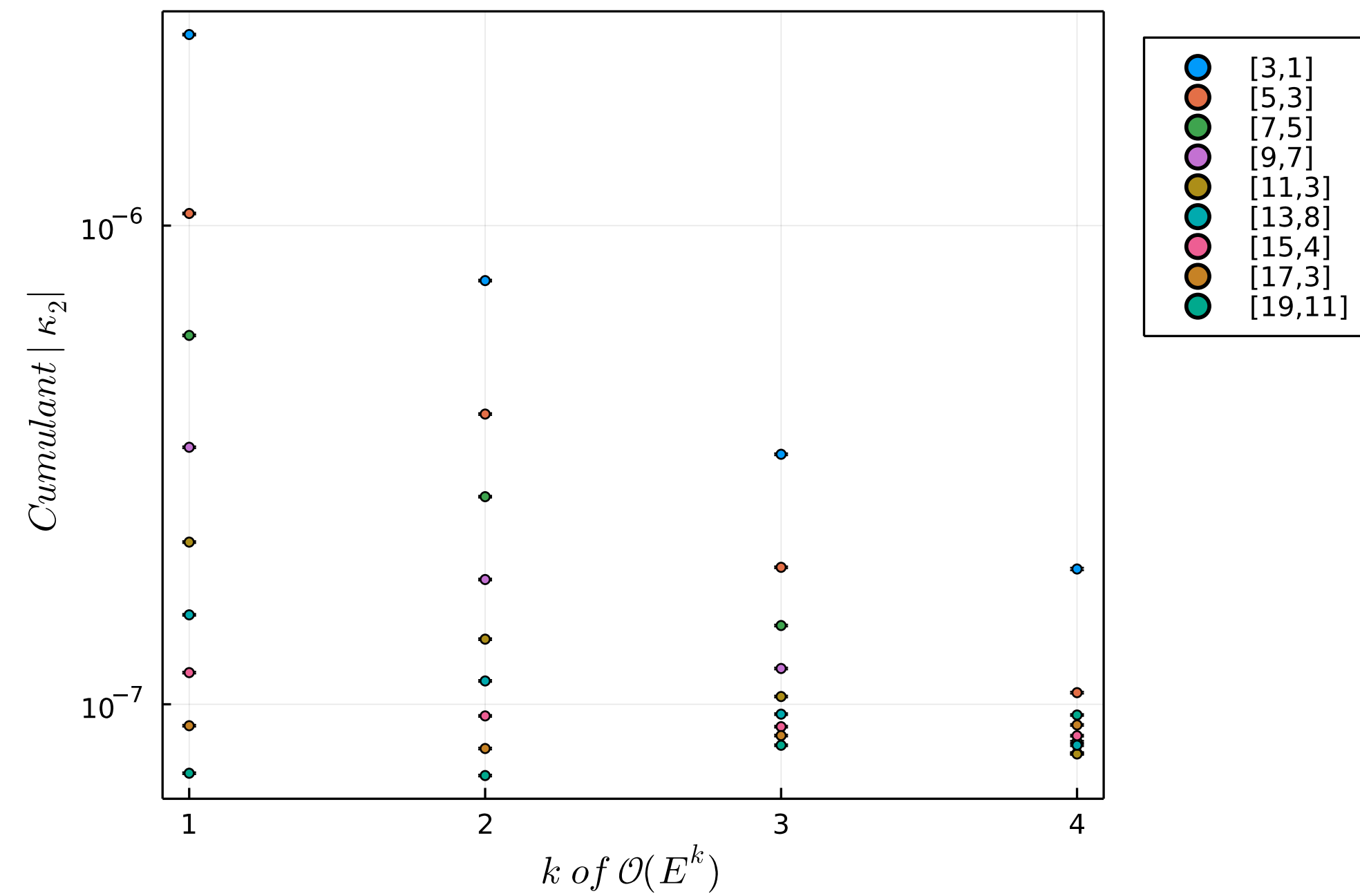


Numerical results

Factorization

$$\langle \mathcal{O}_1 \mathcal{O}_2 \rangle = \langle \mathcal{O}_1 \rangle \langle \mathcal{O}_2 \rangle + \mathbf{O}(1/N_c^2)$$

$$\kappa_2 = \text{var}(E) = \langle E^2 \rangle - \langle E \rangle^2 \Rightarrow 0 \quad \text{as } N_c \rightarrow \infty$$



Summary

- Calculation up to $\mathcal{O}(g^8)$ show the feasibility of combining NSPT and TRPCM.
- The value of the first three coefficients matches very precisely with its theoretical values in the large-N limit.
- The 2 and 3 loop result shows how the volume dependency was eliminated as $N_c \rightarrow \infty$
- 4 loop coefficient with considerable precision.

ENDED

THANK YOU FOR YOUR
WATCHING